

# Impact of nuclear structure on production of superheavy nuclei

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The calculations performed with the modified two-center shell model reveal quite strong shell effects at  $Z = 120 - 126$  and  $N = 184$  [1]. So, our microscopic-macroscopic treatment qualitatively leads to the results close to those in the self-consistent mean-field treatments. If our prediction of the structure of heaviest nuclei is correct, then one can expect the production of evaporation residues  $Z = 120$  in the reactions  $^{50}\text{Ti} + ^{249}\text{Cf}$  and  $^{54}\text{Cr} + ^{248}\text{Cm}$  with the cross sections 23 and 10 fb, respectively. The  $Z = 120$  nuclei with  $N = 175 - 179$  are expected to have  $Q_\alpha$  about 12.1–11.2 MeV and lifetimes 1.7 ms–0.16 s in accordance with our predictions. These  $Q_\alpha$  are in fair agreement with predictions of Liran et al. and about 2 MeV smaller than in other microscopic-macroscopic approaches. The experimental measurement of  $Q_\alpha$  for at least one isotope of  $Z = 120$  nucleus would help us to set proper shell model for the superheavies with  $Z > 118$ .

The nuclear level densities were calculated for the nuclei of  $\alpha$ -decay chains containing  $^{296,298,300}120$ . The minima of the level density parameter clearly indicate the strong shell effect at  $Z = 120$ .

Based on the calculated one-quasiproton spectra and energies for  $\alpha$  decays, one can explain [2] why the  $\alpha$ -decay chain of  $^{291}117$  or  $^{287}115$  is terminated by spontaneous fission of  $^{267}\text{Db}$ . It is shown that, in the  $\alpha$ -decay chain of  $^{293}117$ , the  $\alpha$  decay of  $^{281}\text{Rg}$  is hindered by the structure effects and because of this, the  $^{281}\text{Rg}$  nucleus undergoes spontaneous fission instead of a decay. In addition, the number of isomeric states in the heaviest odd- $Z$  nuclei is predicted. For the  $^{282}\text{Rg}$  nucleus, we expect  $T_\alpha \approx 25$  s and  $T_{sf} \approx 110$  s, i.e., about ten times larger than for neighboring even-odd nuclei  $^{283}\text{Cn}$  and  $^{281}\text{Ds}$ . Thus,  $^{282}\text{Rg}$  undergoes  $\alpha$  decay.

Although the values of  $T_\alpha$  found in the  $\alpha$ -decay chains of  $^{287}\text{Fl}$  and  $^{293}\text{Lv}$  are quite large, one cannot completely exclude the  $\alpha$  decays from the one-quasiparticle isomeric states in  $^{287}\text{Fl}$ ,  $^{283}\text{Cn}$ , and  $^{281}\text{Ds}$  [3]. For example, it is shown that the  $\alpha$  decay of the  $^{281}\text{Ds}$  nucleus occurs only from the ground state which can be populated with small probability in the  $\alpha$ -decay chains of the  $^{289}\text{Fl}$  element. The minimum of  $Q_\alpha$  in  $^{286}\text{Fl}$  indicates the neutron shell at  $N = 172$ .

[1] A.N. Kuzmina, G.G. Adamian, N.V. Antonenko, and W. Scheid, Phys. Rev. C 85, 014319 (2012).

[2] A.N. Kuzmina, G.G. Adamian, and N.V. Antonenko, Phys. Rev. C 85, 017302 (2012).

[3] A.N. Kuzmina, G.G. Adamian, and N.V. Antonenko, Phys. Rev. C 85, 027308 (2012).